

Aaron Z. Goldberg and Anaelle Hertz

National Research Council of Canada

The NRC headquarters is located on the traditional, unceded territory of the Algonquin Anishinaabe and Mohawk peoples.

Coordination or ambiguity



Agenda



Quantifying coherence

From classical to quantum



Quadrature coherence

Many known properties



Spin coherence

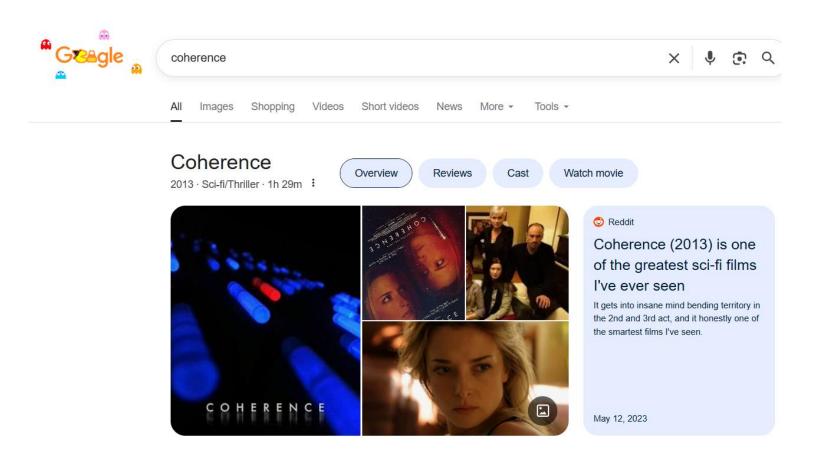
- Desiderata
- Catharsis





The team at NRC (and growing!)





Dictionary

Definition

Synonyms

Example Sentences

Word History

Rhymes

Entries Near

Show More V

Save Word 📑

coherence noun

co·her·ence (kō-ˈhir-ən(t)s ◄»)

-'her- **◄**)

Synonyms of coherence >

- : the quality or state of cohering: such as
 - **a**: systematic or logical connection or consistency

The essay as a whole lacks coherence.

b: integration of diverse elements, relationships, or values

"The various parts of this house—discrete in color, in shape, in placement—join together with remarkable coherence."

- Paul Goldberger

: the property of being coherent

a plan that lacks coherence

Dictionary Definition Synonyms **Example Sentences Word History** Rhymes **Entries Near** Show More V Save Word 💢

coherent adjective

```
co·her·ent (kō-ˈhir-ənt ◄) -ˈher
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Synonyms of *coherent* >

1 a: logically or aesthetically ordered or integrated: CONSISTENT

coherent style

a coherent argument

b: having clarity or intelligibility: **UNDERSTANDABLE**

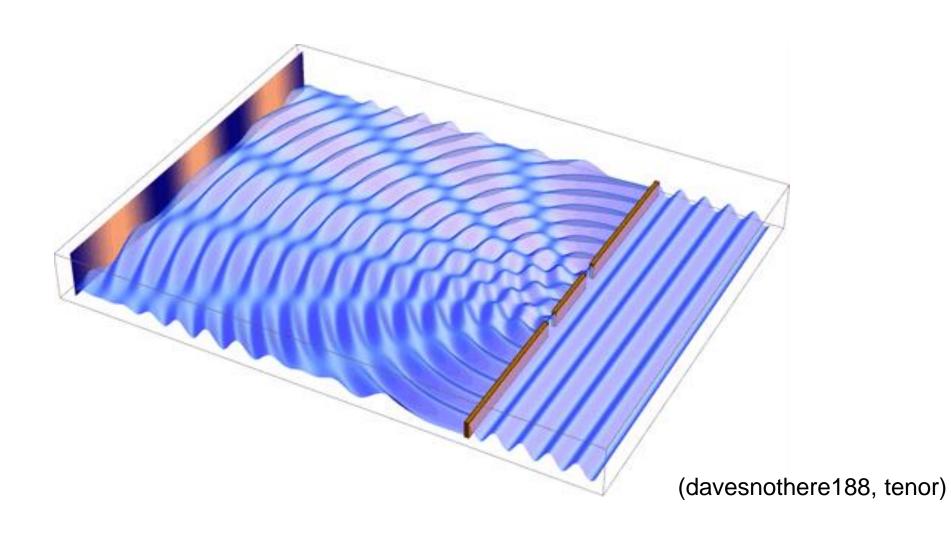
a *coherent* person

a coherent passage

2 : having the quality of holding together or cohering

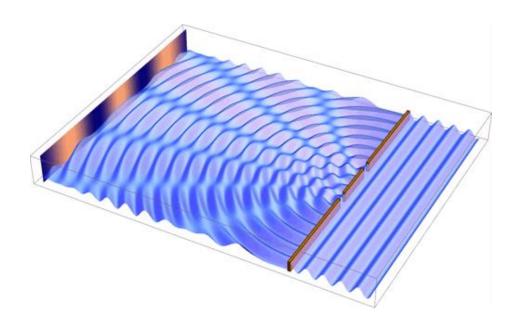
especially: COHESIVE, COORDINATED

a *coherent* plan for action



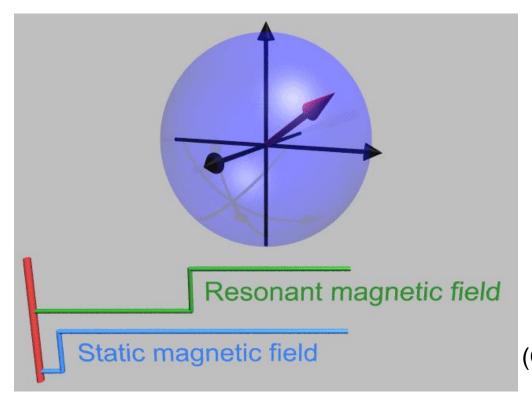
$$Re{\exp[ik(x - x_1)] + \exp[ik(x - x_2)]} = \cdots \cos k(x_1 - x_2)$$

• The ability of two waves to interfere



Quantum: everything is a wave!

Coherent spins will precess in a magnetic field, incoherent won't



(Gavin W Morley, Wikimedia)

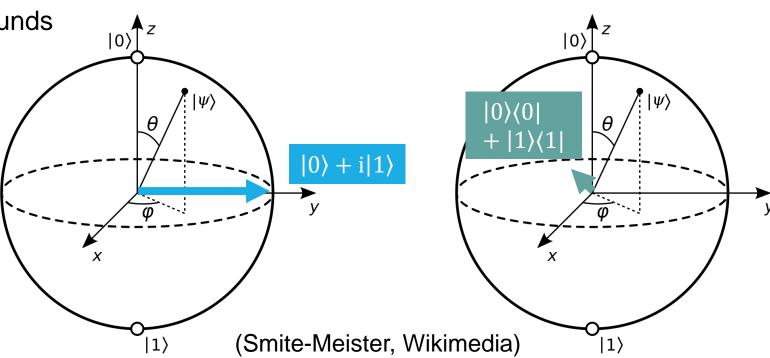
- Optics: compare $\langle E^*(x_1)E(x_2)\rangle$ to $\langle |E(x_1)|^2\rangle$, $\langle |E(x_2)|^2\rangle$
 - Spatial, temporal correlations
- Quantum optics: correlations beyond intensity

Violate classical bounds

• Spins: can precess

• More $|0\rangle + e^{i\phi}|1\rangle$

• Less $|0\rangle\langle 0| + |1\rangle\langle 1|$



For a qubit things are easy

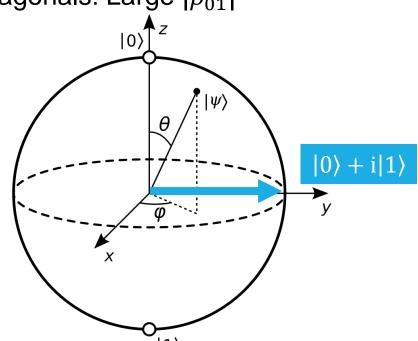
$$\begin{pmatrix} 1 & e^{i\phi} \\ e^{-i\phi} & 1 \end{pmatrix} / 2 \text{ versus } \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} / 2$$

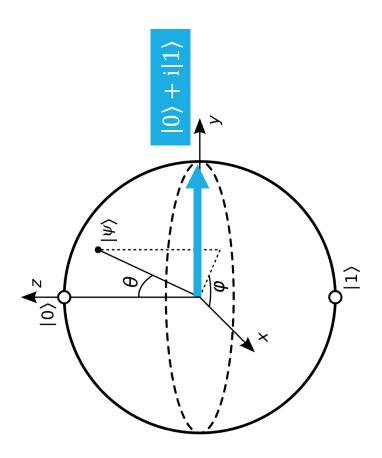
• Magnitude of off-diagonals! Large $|\rho_{01}|$

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- Magnitude of off-diagonals! Large $|\rho_{01}|$
- Not so easy
 - Basis dependent

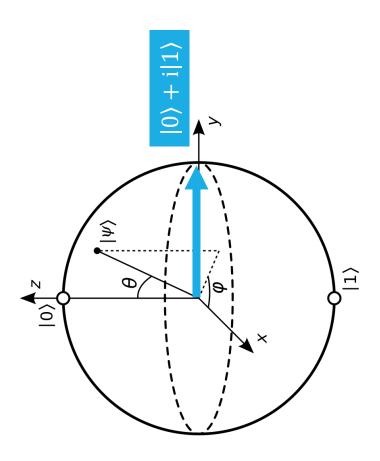




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- Magnitude of off-diagonals! Large $|\rho_{01}|$
- Not so easy
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- Goals for a coherence measure
 - Quantify interferability via off diagonalness
 - Coordinate-system agnostic
 - More than two dimensions



- Single bosonic mode
 - Phase space is position and momentum "quadratures"
 - Infinite dimensions → "continuous variables"
 - Heisenberg-Weyl group generated by x and p with [x, p] = i
 - Superposition of different positions or momenta is uniquely quantum
 - $|x_1\rangle + |x_2\rangle$

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- Somehow, that's easy to calculate

$$C^{2} = -\frac{\text{Tr}([\rho, x]^{2}) + \text{Tr}([\rho, p]^{2})}{2\text{Tr}(\rho^{2})}$$

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- Or divergence of Wigner function
- Or rate of change of purity with loss
- And:
 - Witnesses nonclassicality
 - Bounds distances to classical states
 - Measures quality of displacement sensors (QFI)
 - Convex with loss probably?

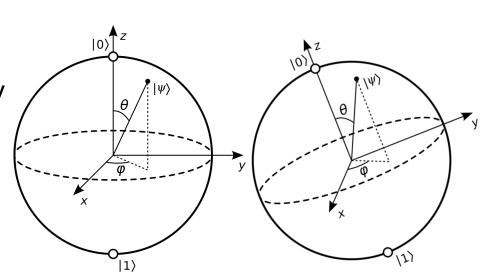
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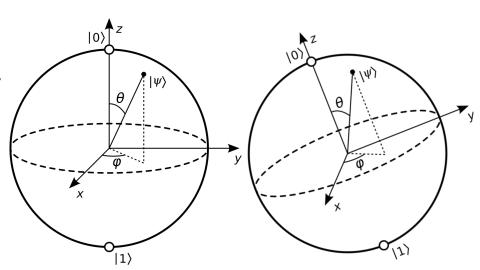
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- And:
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Can we do the same for spins in any finite dimension?

- From Heisenberg-Weyl to SU(2)
 - Three generators; angular momentum $[J_1, J_2] = iJ_3$
 - Casimir invariant or total spin $J_1^2 + J_2^2 + J_3^2 = J(J+1)$ (qubit has J=1/2)
 - Phase space is a sphere, not a plane
 - Get around by rotations $R(\theta, \mathbf{n}) = \exp(i\theta \mathbf{J} \cdot \mathbf{n})$
- Physics: magnetometry, polarimetry, geodesy



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- Physics: magnetometry, polarimetry, geodesy
- Coherence in what basis?
 - Eigenstates $J_3|J,m\rangle = m|J,m\rangle$
 - Or rotate to J_i eigenbasis?
- Noncommutativity of a state $[\rho, J_i]^2$?

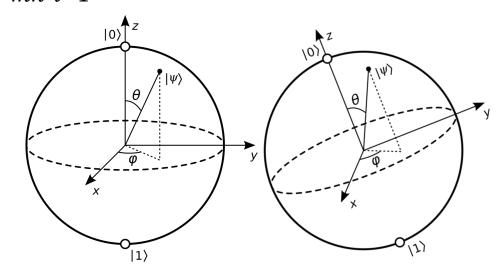


Copy the kid's homework

$$\mathcal{A}^2 \propto \sum_{i=1}^{3} \operatorname{Tr}([\rho, J_i]^2)$$

- 1. Are they equal?
- 2. Basis independent?
- 3. Quantum/classical bound?
- 4. Nice for pure states? Metrology?
- 5. Relations to noise? Monotonicity?
- 6. Quasiprobability distributions?

$$\mathcal{A}^2 \propto \sum_{mn} \sum_{i=1}^{3} (m-n)^2 |_i \langle J, m | \rho | J, n \rangle_i |^2$$



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$$\mathcal{A}^2 \propto \sum_{i=1}^{3} \operatorname{Tr}([\rho, J_i]^2)$$

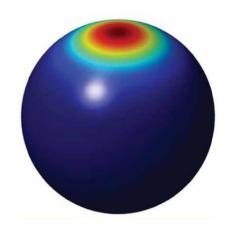
$$\mathcal{A}^2 \propto \sum_{mn} \sum_{i=1}^{3} (m-n)^2 |_i \langle J, m | \rho | J, n \rangle_i |^2$$

- 1. Are they equal? Yes.
- 2. Basis independent? Yes.
- 3. Quantum/classical bound? Yes.
- 4. Nice for pure states? Yes. Metrology? Rotations.
- 5. Relations to noise? Depolarization susceptibility. Monotonicity? Yes.
- 6. Quasiprobability distributions? You bet. Especially for large J

$$\mathcal{A}^2 \propto \sum_{i=1}^3 \operatorname{Tr}([\rho, J_i]^2) \qquad = \qquad \mathcal{A}^2 \propto \sum_{mn} \sum_{i=1}^3 (m-n)^2 |_i \langle J, m | \rho | J, n \rangle_i |^2$$

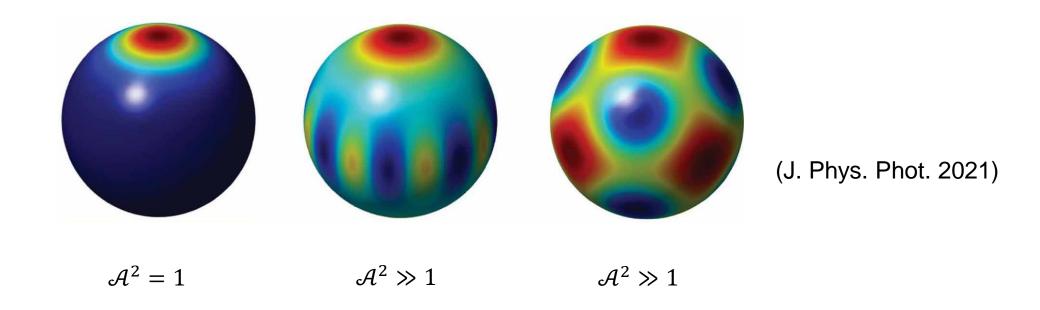
- 1. Are they equal? Yes.
 - Each basis resolves identity $\sum_{m=-J}^{J} |J,m\rangle_i \langle J,m| = \mathbb{I}$
 - Operator ordering sensitivity is the same as coherence scale
- 2. Basis independent? **Yes.** The vector **J** rotates like a vector

- 3. Quantum/classical bound? Yes. What are the classical states?
- Spin coherent states! (Why we refrain from "SCS")
 - Most localized in phase space
 - Minimize $\Delta^2 J_1 + \Delta^2 J_2 + \Delta^3 J_3$, minimize $\Delta^2 J_i \Delta^2 J_j$
 - Equivalent to 2J spin- $\frac{1}{2}$ s all aligned
 - $\langle \Omega | J | \Omega \rangle = J n_{\Omega}$
- Mixtures of coherent states
- Prove $\mathcal{A}^2(\sum_k p_k |\Omega_k\rangle\langle\Omega_k|) \leq 1$ for any probability distribution p_k
 - Not true if $p_k < 0$
 - Proven by $\sum_{i=1}^{3} |\langle \mathbf{\Omega}_k | J_i | \mathbf{\Omega}_l \rangle|^2 \ge J^2 |\langle \mathbf{\Omega}_k | \mathbf{\Omega}_l \rangle|^2$



(J. Phys. Phot. 2021)

- 3. Quantum/classical bound? Yes. What are the classical states?
- Also bounds distance to set of classical states from both above and below



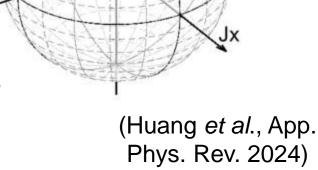
Nice for pure states? Yes. Metrology? Rotations.

$$\mathcal{A}^{2}(|\psi\rangle\langle\psi|) = \frac{1}{J} \sum_{i=1}^{3} \Delta^{2} J_{i}$$

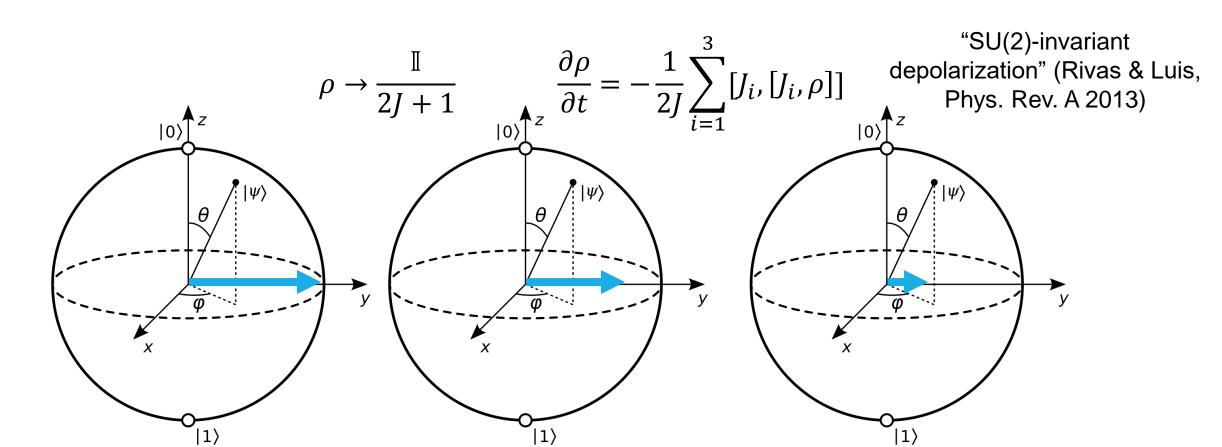
Quantum Fisher information for estimating rotation angle,

averaged over all axes is
$$\frac{1}{4\pi} \int_0^{2\pi} \int_0^{\pi} \text{QFI}(\theta) \ d\Theta \sin \Theta \, d\Phi = \frac{4J}{3} \mathcal{A}^2$$

• Average uncertainty for rotation angle
$$\theta$$
, averaged over all axes
$$\frac{1}{4\pi} \int_0^{2\pi} \int_0^{\pi} \mathrm{Var}(\theta) \ d\Theta \sin\Theta \, d\Phi \geq \frac{3}{4J\mathcal{A}^2}$$



5. Relations to noise? **Depolarization susceptibility.** Monotonicity? **Yes.**

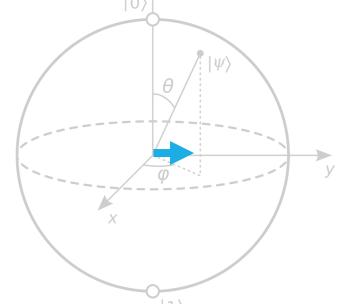


5. Relations to noise? **Depolarization susceptibility.** Monotonicity? **Yes.**

$$\rho \to \frac{\mathbb{I}}{2J+1} \qquad \frac{\partial \rho}{\partial t} = -\frac{1}{2J} \sum_{i=1}^{3} [J_i, [J_i, \rho]]$$

"SU(2)-invariant depolarization" (Rivas & Luis, Phys. Rev. A 2013)

- Then $\mathcal{A}^2 = -\frac{1}{2} \frac{\partial \ln \operatorname{Tr}(\rho^2)}{\partial t}$
- Decreases monotonically with time/noise



- 6. Quasiprobability distributions? You bet. Especially for large J
- s-ordered quasiprobability distributions for SU(2)
 - On the sphere
 - From multipoles, spherical harmonics, Clebsch-Gordan coefficients
- Prove for large *J*:

$$W_{\rho(t)}^{(s)}(\mathbf{\Omega}) \approx W_{\rho(0)}^{(s-4t)}(\mathbf{\Omega})$$

• Just like loss susceptibility for Heisenberg-Weyl $W_{\rho(\eta)}^{(s)}(\alpha) = \frac{1}{\eta} W_{\rho(1)}^{\left(1 - \frac{s-1}{\eta}\right)}(\alpha/\sqrt{\eta})$

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- Just like loss susceptibility for Heisenberg Weyl $W_{\rho(\eta)}^{(s)}(\alpha) = \frac{1}{\eta} W_{\rho(1)}^{\left(1 \frac{s-1}{\eta}\right)}(\alpha/\sqrt{\eta})$
- Then $\mathcal{A}^2 \approx 4 \int W_{\rho}^{(-s)}(\mathbf{\Omega}) \frac{\partial}{\partial s} W_{\rho}^{(s)}(\mathbf{\Omega}) d\mathbf{\Omega} / \int W_{\rho}^{(-s)}(\mathbf{\Omega}') W_{\rho}^{(s)}(\mathbf{\Omega}') d\mathbf{\Omega}'$

Is there more?

- It works for more than spins!
- SU(n), at least in the fundamental representation
 - *n*-mode linear optics
 - Symmetric combinations of *n*-level systems
 - E.g., boson sampling, three-dimensional polarization, polarimetry

Spin coherence scale: operator-ordering sensitivity beyond Heisenberg-Weyl

Aaron Z. Goldberg, Anaelle Hertz

We introduce the spin coherence scale as a measure of quantum coherence for spin systems, generalizing the quadrature coherence scale (QCS) previously defined for quadrature observables. This SU(2)-invariant measure quantifies the off-diagonal coherences of a quantum state in angular momentum bases, weighted by the classical distinguishability of the superposed states. It serves as a witness of nonclassicality and provides both upper and lower bounds on the Hilbert-Schmidt distance to the set of classical (spin coherent) states. We demonstrate that many hallmark properties of the QCS carry over to the spin setting, including its links to noise susceptibility of a state and moments of quasiprobability distributions. The spin coherence scale has direct implications for quantum metrology in the guise of rotation sensing. We also generalize the framework to SU(n) systems, identifying the unique SU(n)-invariant depolarization channel and outlining a broad, Lie-algebraic approach to defining and characterizing the properties of coherence scale beyond harmonic oscillators.

Comments: 12+3 pages; comments always welcome Subjects: Quantum Physics (quant-ph)

Cite as: arXiv:2510.09747 [quant-ph]

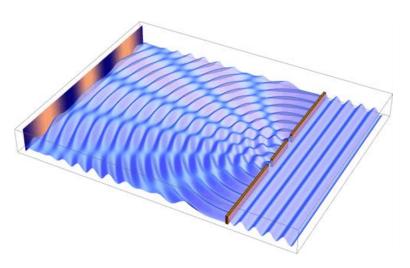
(or arXiv:2510.09747v1 [quant-ph] for this version) https://doi.org/10.48550/arXiv.2510.09747

Conclusions

Coherence is about ambiguous properties being simultaneously present

$$\mathcal{A}^2 \propto \sum_{i=1}^{3} \operatorname{Tr}([\rho, J_i]^2) \propto \sum_{mn} \sum_{i=1}^{3} (m-n)^2 |_i \langle J, m | \rho | J, n \rangle_i |^2$$

- Operator ordering sensitivity, macroscopic superpositionness have ties to...
 - Metrology
 - Quantumness
 - Basis-independent coherence measures
 - Noise susceptibility
 - Quasiprobabilities
- ...beyond Heisenberg-Weyl! All of these properties are

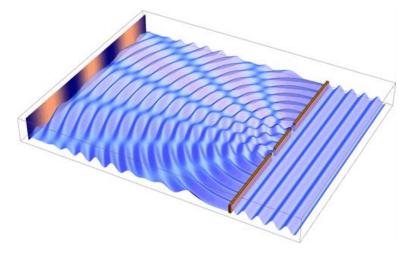


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- ...beyond Heisenberg-Weyl! All of these properties are



coherent.

$$W_{\rho}^{(s)}(\Omega) = \sqrt{\frac{4\pi}{2J+1}} \sum_{Kq} \left(C_{JJ,K0}^{JJ} \right)^{-s} \rho_{Kq} Y_{Kq}(\Omega). \quad (5.17)$$

As is clear, the Clebsch-Gordan coefficients are responsible for the transition between the Husimi function

$$W_{\rho}^{(-1)}(\Omega) = \langle \Omega^{(J)} | \rho | \Omega^{(J)} \rangle \tag{5.18}$$

$$\rho = \sum_{K=0}^{2J} \sum_{q=-K}^{K} \rho_{Kq} T_{Kq}, \tag{5.9}$$

where

$$T_{Kq} = \sqrt{\frac{2K+1}{2J+1}} \sum_{mm'=-J}^{J} C_{Jm,Kq}^{Jm'} |Jm'\rangle\langle Jm|$$
 (5.10)

$$\int d\Omega W_{\rho(0)}^{(s-4t)}(\Omega) W_{\rho(0)}^{(-s+4t)}(\Omega)$$

$$\geq \int d\Omega W_{\rho(0)}^{(s-4t')}(\Omega) W_{\rho(0)}^{(-s+4t')}(\Omega)$$

 $t \le t'$, where $\rho(0)$ can be any state.